Microlocal Methods in Tensor Tomography

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Based on a joint work with GUNTHER UHLMANN

Other collaborators: Bela Frigyik, Venky Krishnan, Nurlan Dirbekov, and Gabriel Paternain

Let (M,g) be a compact Riemannian manifold with boundary.

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If m > 0, you cannot. For any potential field dv, where $v|_{\partial M} = 0$, one has

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The natural conjecture is that this is the only obstruction to uniqueness (for some class of manifolds). We call this property *s-injectivity*.

If m=0 (f is a function), this is just injectivity. By the way, each function α has the same ray transform as αg_{ij} , where g is the metric, so any result for 2-tensors implies a result for functions as well.

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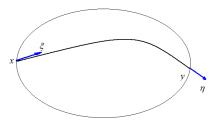
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The Lens Rigidity Problem (The Inverse Kinematic Problem)

Define the scattering relation σ and the length (travel time) function ℓ :



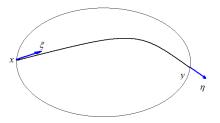
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Do σ , ℓ determine uniquely g, up to an isometry?



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It turns out that linearizing any of those two problems, we arrive at the problem of inverting If modulo potential tensors. Potential tensors linearize the non-uniqueness due to diffeomorphisms.

If m = 0, 1 (f is a function/1-form), then I is (s)-injective on simple manifolds (Mukhometov; Mukhometov & Romanov, Bernstein & Gerver).

If $m \ge 2$, this is still an open problem if $n \ge 3$. In 2D, solved by Sharafutdinov. The non-linear problem was solved (in 2D) before that by PESTOV & UHLMANN.

Energy Estimates

SHARAFUTDINOV, PESTOV: Under an explicit upper bound on the curvature (implying simplicity), I is s-injective with a (non-sharp) stability estimate. DAIRBEKOV: a bit larger class of simple metrics. The energy method goes back to the original idea of MUKHOMETOV but it is a very non-trivial implementation of it on tensors.

Microlocal Approach

S&UHLMANN: Study I^*I as a Ψ DO, and get the most of it. The operator I is an FIO by itself, and this is also used in the analysis. If g is real analytic, use analytic microlocal analysis.

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Every tensor admits an orthogonal decomposition into a solenoidal part f^s and a potential part dv,

$$f = f^s + dv$$
, $v|_{\partial M} = 0$.

where $\delta f^s = 0$.

Here the symmetric differential dv is given by $[dv]_{ij} = (\nabla_i v_j + \nabla_j v_i)/2$, and the divergence δ is given by: $[\delta f]_i = g^{jk} \nabla_k f_{ij}$. We have I(dv) = 0.

To do this, we solve the elliptic boundary value problem

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Note that (1) is sharp, because N is a Ψ DO of order -1.



Idea of the Proof

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One can construct explicitly a parametrix to N in M_1 , more precisely, we "recover" $f_{M_1}^s$ the solenoidal projection of f extended as zero.

We can always assume $f = f^s$. Next, we compare f^s and $f^s_{M_1}$. They differ by some dw, that is known in $M_1 \setminus M$ (up to compact terms) from the parametrix. Then we recover $w|_{\partial M}$ that helps us find f^s .

After that, one gets a Fredholm equation of the kind

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S-injectivity for analytic simple metrics

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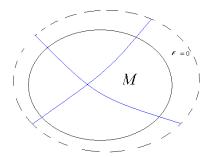
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This is actually an oversimplification of what we are doing. We have to work in M_1 first, that gives us a different solenoidal projection $f_{M_1}^s$.

As an easy example, here is how one can prove this theorem for integrals of functions.

Note that this is a partial case: if f(x) is a function, not a tensor, then $f(x)g_{ij}$ is a tensor, and

$$\int f(\gamma)g_{ij}\dot{\gamma}^i\dot{\gamma}^j\,dt=\int f(\gamma)\,dt.$$

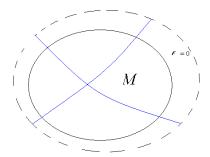


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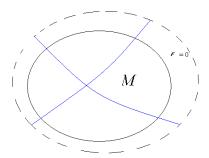


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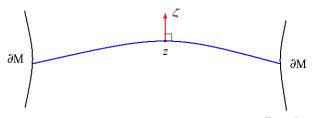
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- incomplete data

Main Condition:

We study $I = I_{\mathcal{D}}$ restricted to $\gamma_{x,\xi}$, where $(x,\xi) \in \mathcal{D} \subset \partial(SM)$. Here \mathcal{D} is chosen so that the conormal bundle of the geodesics issued from \mathcal{D} covers T^*M , and those geodesics have no conjugate points. Such \mathcal{D} are called *complete*.



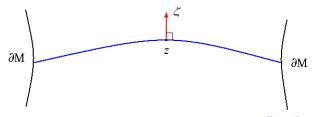
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Definition 3

We say that \mathcal{D} is **complete** for the metric g, if for any $(z,\zeta) \in T^*M$ there exists a maximal in M, finite length geodesic $\gamma: [0,I] \to M$ through z, normal to ζ , such that

- ullet γ belongs to our data (issued from \mathcal{D});
- there are no conjugate points on γ .

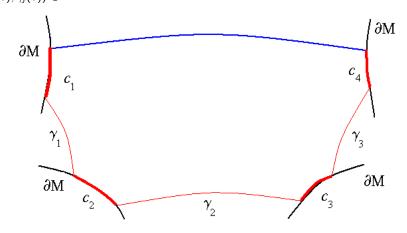
We call g regular, if a complete set \mathcal{D} exists, i.e., if the maximal \mathcal{D} is complete.



Topological Condition (T): Any path in M connecting two boundary points is homotopic to a polygon $c_1 \cup \gamma_1 \cup c_2 \cup \gamma_2 \cup \cdots \cup \gamma_k \cup c_{k+1}$ with the properties that for any j,

(i) c_j is a path on ∂M ;

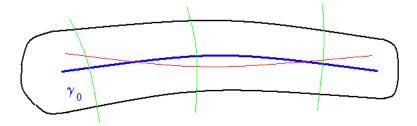
(ii) $\gamma_j:[0,l_j]\to M$ is a geodesic lying in M^{int} with the exception of its endpoints and is transversal to ∂M at both ends; moreover, $(\gamma_i(0),\dot{\gamma}_i(0))\in\mathcal{D}$.



Example 1: A cylinder around an arbitrary geodesic

 γ_0 : a finite length geodesic segment on a Riemannian manifold, conjugate points are allowed.

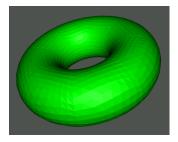
M: a "cylinder" around γ_0 , close enough to it.



One can study the scattering relation only for geodesics almost perpendicular to γ_0 , there are no conjugate points on them.

Example 2: The interior of a perturbed torus

 $M = S^1 \times \{x_1^2 + x_2^2 \le 1\}$, with g close to the flat one:



We need only geodesics almost perpendicular to the boundary. Note that M is trapping!

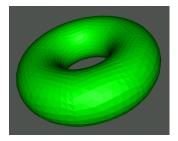
More generally, one can consider a tubular neighborhood of any periodic geodesic on any Riemannian manifold.



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More generally, one can consider a tubular neighborhood of any periodic geodesic on any Riemannian manifold.



Even more generally, we can study $M \times N$, where M is simple, and N is arbitrary; and study σ for all geodesics over fixed points of N, and all those close to them. A small enough perturbation of this manifold satisfies our assumptions, and can have a terrible topology and all kinds of trapping rays and conjugate points.

The examples above are of that type.

Let g be real analytic. Let $\mathcal D$ be open and complete. Then $I_{\mathcal D}$ is s-injective.

Theorem 5 (s-injectivity ⇒ stability)

Let $\mathcal D$ be open and complete. Then s-injectivity of $I_{g,\mathcal D}$ implies a locally uniform stability estimate.

In other words, injectivity implies stability!

Theorem 6 (generic s-injectivity)



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The analytic microlocal arguments in this case

The analytic arguments in this case are much more delicate. First, I^*I may not be a ΨDO at all because we may have conjugate points or trapped geodesics. One can consider

$$N_{\chi} = I^* \chi I$$

$$\chi If = 0 \implies \mathsf{WF}(f) \cap \mathsf{N}^*(\gamma_0) = \emptyset$$



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where χ cuts near some subset of geodesics without conjugate points; for example, near a single one γ_0 . The C^{∞} Ψ DO calculus immediately implies

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If g is analytic, and we want to use analytic ΨDOs , we have a major problem: χ destroys the analyticity. In the analytic ΨDO theory, only certain cut-offs, denoted by $g^R(\xi)$ in Treves' book, are allowed. Here, I is an FIO, and the cut-off is of the type $\chi(x,y,\theta)$, $\theta=\theta(x,y,\xi)$. There is no such theory developed yet.



Instead of trying to use the analytic ΨDO calculus, we use the complex stationary phase method, following Sjöstrand. We get

analytic microlocal arguments

Lemma 7

 γ_0 : non-trapping, without conjugate points. If If = 0 near γ_0 , then

$$WF_{\mathrm{A}}(f^s) \cap N^*(\gamma_0) = \emptyset$$

This is a non-trivial results even for functions

Here is an informal version of those 3 theorems: Under the microlocal and the topologica condition, we have generic s-injectivity for all metrics satisfying those conditions (ncluding the real analytic ones), and a locally uniform stability estimate.

Moreover, the problem is Fredholm (therefore: finite dimensional and smooth kernel).

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Integral geometry of functions over general families of curves.

Consider the weighted X-ray transform of functions over a general family of curves Γ :

$$If(\gamma) = \int w(\gamma(t), \dot{\gamma}(t)) f(\gamma(t)) dt, \quad \gamma \in \Gamma.$$

On can assume that Γ are the solutions of a Newton-type equation

$$\ddot{x} = G(x, \dot{x})$$

with a generator G. (For example, G = 0 gives us lines).

Theorem 8 (Frigyik, S & Uhlmann)

I is injective for generic regular (G, w), including real analytic ones. There is a stability estimate.

Here, G is called regular, if the corresponding curves have no "conjugate points" on supp w, and their conormal bundle (on supp w) covers T^*M . This is the same microlocal condition that we had before, and in particular, we can have a subset of "geodesics".

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Magnetic Systems

On (M,g), consider an one form α , and the Hamiltonian

$$H = \frac{1}{2}(\xi + \alpha)_g^2.$$

The corresponding characteristics on the energy level H=1/2 are called unit speed magnetic geodesics. They describe the trajectories of a charged particle in a magnetic field.

The lens rigidity is formulated in a similar way. The boundary rigidity is formulated in terms of the action A(x,y), on $\partial M \times \partial M$, not the boundary distance function $\rho(x,y)$. The action A(x,y) is defined by

$$A(x,y) = T(x,y) - \int_{\gamma_{[x,y]}} \alpha,$$

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In a joint work with DAIRBEKOV, PATERNAIN AND UHLMANN, we study simple magnetic systems. We prove analogs of the results above. The linearized problem then reduces to the invertibility of the integral transform

$$I\phi(\gamma)=\int_{\gamma}\phi(\gamma,\dot{\gamma})\,dt$$

for functions $\phi(x,\xi)$ that are quadratic in ξ :

$$\phi(x,\xi) = h_{ij}(x)\xi^i\xi^j + \beta_j(x)\xi^j$$

Then *I* is called s-injective, if $I\phi=0$ implies h=dv, $\beta=d\phi-Y(v)$, where $Y(\eta)=((d\alpha)^j_i\eta_i)$.

The uniqueness of the non-linear problem is possible up to a gauge transformation only

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Theorem 9 (S.& Krishnan)

(M,g) simple analytic, K closed geodesically convex subset. If for a symmetric 2-tensor field f we have that $If(\gamma) = 0$ for each geodesic γ not intersecting K, then there exists an 1-form v such that f = dv in $M \setminus K$, and v = 0 on ∂M .

True for functions (Venky Krishnan) and 1-forms as well, and the proof if simple:

Working with f^s is not what we should do now. It is not true in general that $f^s=0$ in $M\setminus K$. After adding some dv, one can always assume that

$$f_{ni}=0, \forall i$$

near a fixed geodesic γ_0 , in special coordinates, where $\gamma_0 = (0, \dots, 0, t)$. Now, I (or $N = I^*I$) is not elliptic on such tensors, but it is elliptic for ξ with $\xi_n \neq 0$ (not conormal to γ_0).

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the analytic microlocal argument

One of the ingredients of the proof is a refined version of the analytic microlocal argument discussed above:

Lemma 10

 γ_0 : non-trapping, without conjugate points. Given $(x_0, \xi^0) \in N^* \gamma_0$, assume that $(x_0, \xi^0) \notin WF_A(\delta f)$, and that If = 0 near γ_0 . Then

$$(x_0,\xi^0) \not\in WF_{\mathcal{A}}(f).$$

Another ingredient is the Sato-Kawai-Kawashita Theorem: Let f be supported on one side of a hypersurface S, and $x_0 \in S$, $\xi^0 \perp S$ at x_0 . Assume that f is analytic at (x_0, ξ^0) . Then f = 0 near x_0 .

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Modeling Optical Molecular Imaging

The radiative transport equation in Ω is given by

$$\theta \cdot \nabla_{x} u(x,\theta) + \sigma(x,\theta) u(x,\theta) - \int_{S^{n-1}} k(x,\theta,\theta') u(x,\theta') d\theta' = f(x), \quad u|_{\partial_{-}S\Omega} = 0,$$

where σ is the absorption and k is the collision kernel. The source term f is assumed to depend on x only. Here, $\partial_- S\Omega$ consists of $x \in \partial\Omega$ and θ pointing inwards.

The boundary measurements are modeled by

$$Xf(x,\theta) = u|_{\partial_{+}S\Omega}, \quad (x,\theta) \in \partial_{+}S\Omega$$

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Given f (and σ , k), find Xf.

Inverse Problem

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Clearly, it is a linear problem.

Let $\sigma = k = 0$ first

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Let k = 0 only.

Then we get a weighted X-ray transform:

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where

$$E(x,\theta) = \exp\left(-\int_0^\infty \sigma(x+s\theta,\theta)\,ds\right).$$

If $\sigma=\sigma(x)$, then we get the attenuated X-ray transform, that we know how to invert.

Without assuming that any one is zero, BAL and TAMASAN proved injectivity when $k = k(x, \theta \cdot \theta')$, and k is small enough in a suitable norm. The main idea there is to treat k as a perturbation; then X is a perturbation of the attenuated X-ray transform.

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Our goal is to consider this problem for general (σ, k) .



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The direct problem first:

We need assumptions, even for solvability of the direct problem! Assuming $|k| \ll 1$ is enough. Also, $\int k(\cdot, \theta, \cdot) d\theta < \sigma$ suffices. Those conditions prevent a "nuclear explosion", i.e., the corresponding time-dependent dynamics is bounded. They are not necessary conditions though.



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Theorem 11

- (a) The direct problem is uniquely solvable for a dense open set of pairs (σ, k) in C^2 , even for $f = f(x, \theta)$.
- (b) $X: L^2(\Omega \times S^{n-1}) \longrightarrow L^2(\partial_+ S\Omega, d\Sigma)$.



The direct problem first:

We need assumptions, even for solvability of the direct problem! Assuming $|k| \ll 1$ is enough. Also, $\int k(\cdot, \theta, \cdot) d\theta < \sigma$ suffices. Those conditions prevent a "nuclear explosion", i.e., the corresponding time-dependent dynamics is bounded. They are not necessary conditions though.

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Here. $d\Sigma = \nu \cdot \theta dS_{\star} d\theta$.



The Inverse Problem

Fix $\Omega_1 \supset \supset \Omega$. Define X_1 as X but in Ω_1 .

Theorem 12

For (σ, k) in an open and dense set of pairs in

$$C^2(\bar{\Omega}\times S^{n-1})\times C^2\left(\bar{\Omega}_{\mathsf{x}}\times S^{n-1}_{\theta'};\ C^{n+1}(S^{n-1}_{\theta})\right)$$

including (0,0), the direct problem is solvable in Ω_1 , and

- (a) the map X_1 is injective on $L^2(\Omega)$,
- (b) the following stability estimate holds

$$||f||_{L^2(\Omega)} \le C||X_1^*X_1f||_{H^1(\Omega_1)}, \quad \forall f \in L^2(\Omega),$$

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